

Superconductivity from confinement transition of FL* metals with Z_2 topological order

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Z_2 spin liquids

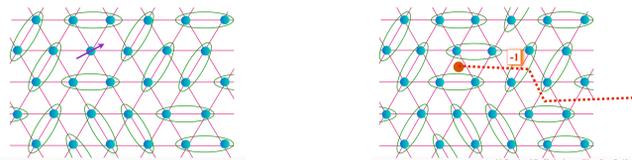
- Quantum disordered ground states of certain Mott insulators
- No broken symmetry, topological degeneracy of ground states
- Schwinger boson or Abrikosov fermion mean-field theory:

$$\vec{S} = \frac{1}{2} b_{i\alpha}^\dagger \vec{\sigma}_{\alpha\beta} b_{i\beta}, \text{ or } \vec{S} = \frac{1}{2} f_{i\alpha}^\dagger \vec{\sigma}_{\alpha\beta} f_{i\beta}$$

Read and Sachdev, PRL **66**, 1773 (1991)
Wen, PRB **44**, 2662 (1991)

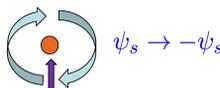
$$H_{MF}^b = - \sum_{i,j} (Q_{ij} \epsilon_{\alpha\beta} b_{i\alpha}^\dagger b_{j\beta}^\dagger + h.c.) + \sum_i \lambda_i (b_{i\alpha}^\dagger b_{i\alpha} - 1)$$

Excitations:



- Fractionalized spin-half spinons
- Visons or vortices of the Z_2 gauge field

Spinons and visons are mutual semions



J_1 - J_2 - J_3 Heisenberg model on the square lattice:

- Quantum fluctuations can drive a continuous phase transition from a spiral incommensurate antiferromagnet at $(Q,0)$ or $(0,Q)$ to a Z_2 spin liquid with Ising nematic order \Rightarrow broken C_4 symmetry $\Rightarrow |Q_{i,i+\hat{x}}|^2 - |Q_{i,i+\hat{y}}|^2 \neq 0$

Linking bosonic and fermionic descriptions

Topological properties of Z_2 gauge theory:

- 4 kinds of topologically distinct excitations: e , m , ϵ and 1 (topologically trivial)

- Fusion rules: $e \times e = m \times m = \epsilon \times \epsilon = 1$
 $1 \times 1 = 1, e \times 1 = e, m \times 1 = m, \epsilon \times 1 = \epsilon$



$$e \times m = \epsilon, e \times \epsilon = m, m \times \epsilon = e \quad \text{Kitaev, Annals of Physics 321 (2006)}$$

In the context of Z_2 spin liquids:

$$\begin{aligned} e &\rightarrow b \text{ (bosonic spinon)} \\ \epsilon &\rightarrow f \text{ (fermionic spinon)} \\ m &\rightarrow v \text{ (vison)} \end{aligned}$$

Symmetry fractionalization:

Essin and Hermele, PRB **87**, 104406 (2013)
Lu, Cho and Vishwanath, arXiv:1403.0575

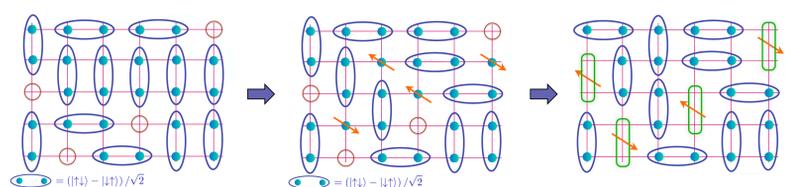
- Symmetries act projectively on individual anyons
- For each symmetry (space-group, time-reversal) combination equivalent to identity, we can associate a Z_2 quantum number for each anyon
- From the Z_2 quantum numbers of the bosonic spinons and visons, we can determine the symmetry fractionalization quantum numbers of the fermionic spinons for a fully gapped Z_2 spin liquid

Equivalence of bosonic and fermionic Z_2 spin liquids on the rectangular lattice:

Symmetry commutation	σ_O^e	σ_O^m	σ_O^f	σ_O^v
$T_x^{-1} T_y^{-1} T_x T_y$	$(-1)^{p_1}$	-1	1	$\sigma_{T_x T_y}^e = (-1)^{p_1+1}$
$P_x^{-1} T_x P_x T_x$	$(-1)^{p_2}$	1	1	$\sigma_{P_x T_x}^e = (-1)^{p_2}$
$P_y^{-1} T_y P_y T_y$	$(-1)^{p_3}$	-1	1	$\sigma_{P_y T_y}^e = (-1)^{p_3+1}$
$P_x^{-1} T_x^{-1} P_x T_x$	$(-1)^{p_4}$	-1	1	$\sigma_{P_x T_x}^e = (-1)^{p_4+1}$
$P_y^{-1} T_y^{-1} P_y T_y$	$(-1)^{p_5}$	1	1	$\sigma_{P_y T_y}^e = (-1)^{p_5}$
P_x^2	$(-1)^{p_6}$	1	-1	$\sigma_{P_x}^e = (-1)^{p_6+1}$
P_y^2	$(-1)^{p_7}$	1	-1	$\sigma_{P_y}^e = (-1)^{p_7+1}$
$T_x^{-1} P_y^{-1} P_x P_y$	1	-1	-1	$\sigma_{P_x P_y}^e = 1$
T^2	-1	1	1	$\sigma_T^e = -1$
$T_x^{-1} T^{-1} T_x T$	$(-1)^{p_8}$	1	1	$\sigma_{T T_x}^e = (-1)^{p_8}$
$T_y^{-1} T^{-1} T_y T$	$(-1)^{p_9}$	1	1	$\sigma_{T T_y}^e = (-1)^{p_9}$
$P_x^{-1} T^{-1} P_x T$	$(-1)^{p_6}$	1	-1	$\sigma_{T P_x}^e = (-1)^{p_6+1}$
$P_y^{-1} T^{-1} P_y T$	$(-1)^{p_7}$	1	-1	$\sigma_{T P_y}^e = (-1)^{p_7+1}$

- Can write down a Hamiltonian consistent with the projective symmetry realization for the fermionic spinons \Rightarrow Pi-flux gapped spin liquid

Fractionalized Fermi liquid (Z_2 FL*)



- Metallic state with charge- e spin-half c fermions in the background of a Z_2 spin liquid

- The size of the Fermi surface is determined by dopant density p

- No low energy fractionalized excitations

- The vortices of the internal Z_2 gauge field (visons) survive in the FL* metal, hence its topological character

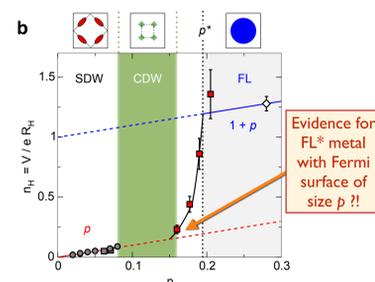
- Violates Luttinger's theorem due to presence of emergent gauge excitations

Additional sectors of Z_2 -FL*

	1_c	ϵ_c	m_c	ϵ_c
S	1/2	0	1/2	0
Statistics	fermion	fermion	fermion	boson
Mutual semions	-	$m, \epsilon, m_c, \epsilon_c$	$\epsilon, \epsilon, \epsilon_c, \epsilon_c$	$\epsilon, m, \epsilon_c, m_c$
Q	1	1	1	1
Field operator	c	-	-	B

bosonic chargon

Hall effect measurements in YBCO



Badoux, Proust, Taillefer et al, Nature **531**, 9552 (2016)

Superconducting transition

Confinement transitions out of the Z_2 FL*:

Condense $b \Rightarrow$ SDW

Condense $v \Rightarrow$ Bond-DW

Condense $B \Rightarrow$ Superconductivity

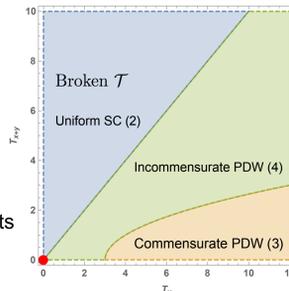
Chubukov, Senthil and Sachdev, PRL **72**, 2089 (1994)

Patel, Chowdhury, Allais and Sachdev, PRB **93**, 165139 (2016)

- A Higgs transition induced by condensation of $B \sim f_{i\alpha}^\dagger c_\alpha$ or $\epsilon_{\alpha\beta} c_{i\alpha} f_{i\beta}$ results in confinement as B carries Z_2 gauge charge

- The pairing of f spinons induces a pairing of the c fermions, resulting in a superconductor

- Non-trivial transformation of B under translation results in spatial modulation of the superconducting order parameter (Pair-Density Wave or FFLO state)



D-wave SC, Density waves and Hall effect

- Consider a 'plain vanilla' Z_2 FL* with trivial PSG for fermionic spinons

- D-wave spinon-pairing leads to uniform d-wave SC

- Modified boson dispersions can lead to co-existing uniform SC with bond density waves and pair density waves at the same axial wave-vector Q , as observed in recent STM experiments. Fujita et al, PNAS **111**, 3026 (2014) Hamidian et al, Nature **532**, 343 (2016)

- Hall No. also shows a jump at the optimal doping critical point in the metallic state at $T=0$

Open questions:

- How does the jump in the Hall No. smoothen at finite T ?

- Which experiments can distinguish the Z_2 FL* from other candidate topological metals or field-induced magnetism?

